

# SENSITIVITY ANALYSIS FOR SZEGŐ POLYNOMIALS

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**Abstract.** Szegő polynomials are orthogonal with respect to an inner product on the unit circle. Numerical methods for weighted least-squares approximation by trigonometric polynomials conveniently can be derived and expressed with the aid of Szegő polynomials. This paper discusses the conditioning of several mappings involving Szegő polynomials and, thereby, sheds light on the sensitivity of some approximation problems involving trigonometric polynomials.

**1. Introduction.** Let  $\mu(t)$  be a distribution function with infinitely many points of increase in the interval  $[-\pi, \pi]$  and let

$$\mu_j := \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-ijt} d\mu(t), \quad j = 0, \pm 1, \pm 2, \dots,$$

be the moments associated with  $\mu(t)$ , where  $i := \sqrt{-1}$ . For notational convenience, we assume throughout this paper that  $\mu(t)$  is scaled so that  $\mu_0 = 1$ . We note that the moment matrices  $M_n := [\mu_{j-k}]_{j,k=0}^{n-1}$ ,  $n = 1, 2, 3, \dots$ , are Hermitian positive definite and of Toeplitz form. There is an infinite sequence of monic polynomials  $\{\psi_j\}_{j=0}^{\infty}$ , known as *Szegő polynomials*, that are orthogonal with respect to the inner product on the unit circle

$$(1.1) \quad (f, g) := \frac{1}{2\pi} \int_{-\pi}^{\pi} \overline{f(e^{it})} g(e^{it}) d\mu(t),$$

where the bar denotes complex conjugation. Thus,  $(\psi_j, \psi_k) = 0$  for  $j \neq k$ , and  $(\psi_j, \psi_j) > 0$  for all  $j \geq 0$ .

Properties of Szegő polynomials are discussed in, e.g., [15, 21, 26, 29]. Of particular importance is the fact that the monic Szegő polynomials satisfy recurrence relations of the form

$$(1.2) \quad \begin{aligned} \psi_j(z) &= z\psi_{j-1}(z) + \gamma_j \tilde{\psi}_{j-1}(z), \\ \tilde{\psi}_j(z) &= \bar{\gamma}_j z\psi_{j-1}(z) + \psi_{j-1}(z), \end{aligned} \quad j = 1, 2, 3, \dots,$$

with  $\psi_0(z) := 1$  and  $\tilde{\psi}_0(z) := 1$ . The functions  $\tilde{\psi}_j$  often are referred to as *reverse polynomials*, because they satisfy  $\tilde{\psi}_j(z) = z^j \bar{\psi}_j(z^{-1})$ ; thus, the coefficients of  $\tilde{\psi}_j(z)$ , when represented in terms of powers of  $z$ , are the conjugated coefficients of  $\psi_j(z)$  in reverse order. The recursion coefficients  $\gamma_j \in \mathbb{C}$ , also known as *Schur parameters*, and the complementary parameters  $\sigma_j \in \mathbb{R}^+$  are determined by

$$\begin{aligned} \gamma_j &= -(1, z\psi_{j-1})/\sigma_{j-1}, \\ \sigma_j &= \sigma_{j-1}(1 - |\gamma_j|^2), \end{aligned} \quad j = 1, 2, 3, \dots,$$

with  $\sigma_0 = 1$ ; see, e.g., [15, 21, 26, 29]. It is well known that  $|\gamma_j| < 1$  for  $j \geq 1$ , and that all the zeros of Szegő polynomials lie in the open unit disk in  $\mathbb{C}$ .

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Consider the integral

$$I(f) := \frac{1}{2\pi} \int_{-\pi}^{\pi} f(e^{it}) d\mu(t).$$

An  $n$ -point Szegő quadrature rule for this integral is of the form

$$(1.3) \quad S_{\mu}^{(n)}(f) := \sum_{j=1}^n \omega_j^2 f(\lambda_j), \quad \omega_j^2 > 0, \lambda_j \in \Gamma,$$

where  $\Gamma$  denotes the unit circle in  $\mathbb{C}$ . The defining property of  $S_{\mu}^{(n)}$  is that

$$(1.4) \quad S_{\mu}^{(n)}(p) = I(p), \quad \forall p \in \Lambda_{-(n-1), n-1},$$

where  $\Lambda_{-(n-1), n-1}$  is the linear space of Laurent polynomials

$$L(z) := \sum_{k=-(n-1)}^{n-1} c_k z^k, \quad c_k \in \mathbb{C},$$

of degree at most  $n-1$ . Equivalently,  $S_{\mu}^{(n)}$  is exact for all trigonometric polynomials of degree less than  $n$ . In particular,

$$(1.5) \quad \sum_{j=1}^n \omega_j^2 = \mu_0 = 1.$$

Differently from the situation for Gauss quadrature rules, the nodes of Szegő rules are not uniquely determined. One node can be chosen arbitrarily on  $\Gamma$ ; see, e.g., [12, 15, 19, 21] for discussions on Szegő quadrature rules.

Following Gragg [12], we express the recursion relations (1.2) for  $1 \leq j \leq n$ , with  $\gamma_n = \tau$ , in matrix form. Define the upper Hessenberg matrix

$$\widehat{H}_{\tau} := \begin{bmatrix} -\bar{\gamma}_0 \gamma_1 & -\bar{\gamma}_0 \gamma_2 & \cdots & -\bar{\gamma}_0 \gamma_{n-1} & -\bar{\gamma}_0 \tau \\ 1 - |\gamma_1|^2 & -\bar{\gamma}_1 \gamma_2 & \cdots & -\bar{\gamma}_1 \gamma_{n-1} & -\bar{\gamma}_1 \tau \\ & 1 - |\gamma_2|^2 & \cdots & -\bar{\gamma}_2 \gamma_{n-1} & -\bar{\gamma}_2 \tau \\ & & \ddots & \vdots & \vdots \\ 0 & & & 1 - |\gamma_{n-1}|^2 & -\bar{\gamma}_{n-1} \tau \end{bmatrix}$$

with  $\gamma_0 := 1$ . Then, for  $1 \leq j \leq n$ , the recursions (1.2) can be written as

$$[\psi_0(z), \psi_1(z), \dots, \psi_{n-1}(z)] \widehat{H}_{\tau} = z[\psi_0(z), \psi_1(z), \dots, \psi_{n-1}(z)] - [0, \dots, 0, \Psi_{n,\tau}(z)],$$

where  $\Psi_{n,\tau}(z)$  is a polynomial of degree  $n$  in  $z$ . Its zeros are the eigenvalues of  $\widehat{H}_{\tau}$ . If  $\tau = \gamma_n$ , then  $\Psi_{n,\tau} = \psi_n$ , the  $n$ th Szegő polynomial, and the eigenvalues of  $\widehat{H}_{\tau}$  are the zeros of this polynomial. Note that  $\widehat{H}_{\tau}$  may have multiple eigenvalues and be defective. This is the case when  $\mu(t) := t$  and  $\tau := 0$ . Then  $\psi_j(z) = z^j$  for  $j = 0, 1, 2, \dots$ .

We will let  $\tau \in \Gamma$ . Gragg [12] showed that for such  $\tau$ , the matrix

$$(1.6) \quad H_{\tau} := D_n^{-1/2} \widehat{H}_{\tau} D_n^{1/2}$$

with  $D_n := \text{diag}[\sigma_0, \sigma_1, \dots, \sigma_{n-1}]$  is unitary, and that its eigenvalues are the nodes,  $\lambda_j$ , and the squared magnitude of the first component of associated normalized eigenvectors are the weights  $\omega_j^2$  of an  $n$ -point Szegő rule (1.3). The rule satisfies (1.4) for all  $\tau \in \Gamma$ . We will denote the matrix  $H_\tau$  simply by  $H$ , since the  $\tau$ -dependence is not important for our purpose. Efficient algorithms for the computation of Szegő quadrature rules are presented in [13, 14, 16, 30, 31].

Szegő polynomials arise in many applications, such as in signal processing, time series analysis, and operator theory; see, e.g., [1, 4, 15, 20, 22, 23, 24]. In particular, least-squares approximation by trigonometric polynomials gives rise to the following *inverse eigenvalue problem* for the unitary upper Hessenberg matrix (1.6): Given  $n$  distinct eigenvalues  $\{\lambda_j\}_{j=1}^n$  on  $\Gamma$  and the first components  $\{\omega_j\}_{j=1}^n$  of the corresponding eigenvectors normalized to have unit length with  $\omega_j > 0$ , determine the unique unitary upper Hessenberg matrix  $H$  with positive subdiagonal elements, such that

$$HW = W\Lambda,$$

where  $\Lambda = \text{diag}[\lambda_1, \lambda_2, \dots, \lambda_n]$  and  $W$  is a unitary eigenvector matrix with

$$(1.7) \quad e_1^t W e_j = \omega_j, \quad 1 \leq j \leq n.$$

Throughout this paper the superscript  $t$  denotes transposition, and the superscript  $*$  transposition and complex conjugation. The uniqueness of the matrix  $H$  follows from the Implicit Q Theorem [11, Theorem 7.4.2]. Existence and uniqueness of the solution of this inverse eigenvalue problem are discussed in [2, 25], where also algorithms for the computation of the solution are presented.

We are interested in investigating the sensitivity of the mapping

$$(1.8) \quad \mathcal{F}_n : [\omega, \lambda] \mapsto H, \quad [\omega, \lambda] := [\omega_1, \omega_2, \dots, \omega_n, \lambda_1, \lambda_2, \dots, \lambda_n]$$

to perturbations in the data and, in particular, how the size of the  $\omega_j$  and the location of the  $\lambda_j$  affects the conditioning of this map.

A fundamental problem in Scientific Computation is the determination of orthogonal polynomials from given moments. Several algorithms for computing Szegő polynomials from the moments have been proposed; see, e.g., [12, 17, 24, 25]. It is therefore of interest to study the conditioning of the mapping

$$\mathcal{K}_n : \mu \mapsto H, \quad \mu := [\mu_0, \mu_1, \dots, \mu_{n-1}].$$

It is convenient to consider  $\mathcal{K}_n$  a composition of two maps  $\mathcal{K}_n = \tilde{\mathcal{F}}_n \circ \mathcal{G}_n$ , where

$$(1.9) \quad \tilde{\mathcal{F}}_n : [\omega^2, \lambda] \mapsto H, \quad [\omega^2, \lambda] := [\omega_1^2, \omega_2^2, \dots, \omega_n^2, \lambda_1, \lambda_2, \dots, \lambda_n],$$

maps the quadrature data to the matrix (1.6) and

$$\mathcal{G}_n : \mu \mapsto [\omega^2, \lambda]$$

maps the moments to the quadrature data. This approach to studying the conditioning of  $\mathcal{K}_n$  has been applied by Gautschi [8, 9] in his investigations of the sensitivity of polynomials orthogonal with respect to an inner product on an interval on the real axis as a function of the moments or modified moments.

Having determined a bound for the condition number of the map  $\mathcal{F}_n$ , we easily can compute a bound for the condition number of  $\tilde{\mathcal{F}}_n$ . This is carried out in Section

2. The conditioning of the mapping  $\mathcal{G}_n$  is investigated in [18]. Define the condition number

$$(1.10) \quad \kappa(\mathcal{G}_n)(\mu) := \frac{\|\mu\|_2}{\|[\omega^2, \lambda]\|_2} \|J_{\mathcal{G}_n}\|_2,$$

where  $J_{\mathcal{G}_n}$  is the Jacobian of  $\mathcal{G}_n$ . Then, it is shown in [18] that

$$(1.11) \quad \kappa(\mathcal{G}_n)(\mu) < \frac{1}{\sqrt{2n-1}} \cdot \left\{ \sum_{k=1}^{2n-1} |P_1(e^{i\theta_k})|^2 + \sum_{j=2}^n \sum_{k=1}^{2n-1} \left( |P_j(e^{i\theta_k})|^2 + \frac{|Q_j(e^{i\theta_k})|^2}{\omega_j^2} \right) \right\}^{1/2},$$

where  $\theta_k = 2\pi k/(2n-1)$  and  $P_j, Q_j$  are the Laurent polynomials

$$\begin{aligned} P_1(z) &= l_1(z)\bar{l}_1(z^{-1}), \\ P_j(z) &= \hat{l}_j(z)\bar{l}_j(z^{-1}) \left\{ \frac{z - \lambda_1}{\lambda_j - \lambda_1} - \left( \hat{l}'_j(\lambda_j) - \lambda_j^{-2}\bar{l}'_j(\lambda_j^{-1}) \right) (z - \lambda_j) \right\}, \\ Q_j(z) &= \hat{l}_j(z)\bar{l}_j(z^{-1})(z - \lambda_j). \end{aligned}$$

Here  $l_j(z)$  and  $\hat{l}_j(z)$  are the Lagrange polynomials defined by

$$(1.12) \quad l_j(z) := \prod_{\substack{k=1 \\ k \neq j}}^n \frac{z - \lambda_k}{\lambda_j - \lambda_k}, \quad \hat{l}_j(z) := \prod_{\substack{k=2 \\ k \neq j}}^n \frac{z - \lambda_k}{\lambda_j - \lambda_k}.$$

The bound (1.11) is modest unless there is a “tiny” weight  $\omega_j^2$ , there are pairs of close nodes  $\lambda_j$ , or most nodes are clustered on a small subarc of  $\Gamma$ . We will see that, except for in these situations, the mapping  $\tilde{\mathcal{F}}_n$  also is fairly well-conditioned. We therefore may conclude that the map  $\mathcal{K}_n$  is fairly well-conditioned when there are no tiny weights and the nodes are fairly evenly spread out over  $\Gamma$  with no pair of adjacent nodes very close.

It is well known that the map from the moments to the recursion coefficients for polynomials orthogonal with respect to an inner product on an interval on the real axis generally is severely ill-conditioned; the condition number typically grows rapidly and exponentially with  $n$ . Therefore, Gautschi studied the use of so-called *modified moments* instead of the ordinary moments, and showed that the determination of recursion coefficients from modified moments can be less sensitive to perturbations in the modified moments; see [7, 8, 9]. The modified moments are moments with respect to a family of polynomials other than the power basis. It is usually advantageous to choose the family to be a sequence of orthogonal polynomials with respect to an inner product on  $\mathbb{R}$ , such as a sequence of Chebyshev or Laguerre polynomials; see Beckermann and Bourreau [3], as well as Fischer [5] and Gautschi [7, 8, 9] for discussions. Related results for Sobolev inner products on a real interval are discussed by Zhang [32]. Gautschi and Zhang [10] describe numerical methods using modified moments in this context.

The monomials  $z^j$  are orthogonal with respect to the inner product (1.1) when  $\mu(t) = t$ . Therefore the moments  $\mu_j$  are analogous to modified moments. This suggests that the determination of recursion coefficients for Szegő polynomials from

moments generally should be less sensitive to perturbations than the computation of recursion coefficients for polynomials orthogonal with respect to an inner product on the real axis from the moments. Combining results of this paper with those in [18] supports this observation. In particular, the numerical examples of Section 3 show the condition numbers of  $\mathcal{F}_n$  and  $\tilde{\mathcal{F}}_n$  to grow slower than exponentially.

The usefulness of Szegő polynomials in signal processing stems from their connection with symmetric positive definite Toeplitz matrices. In signal processing applications the recursion coefficients  $\gamma_j$  often are computed by factoring a Toeplitz autocorrelation matrix associated with a wide sense stationary signal by the Schur or Levinson algorithms; see, e.g., [1, 12, 15, 20, 22, 23]. Block generalizations of the latter algorithms are available and their numerical properties are discussed in, e.g., [6, 28]. These block algorithms compute matrix-valued recursion coefficients, which determine matrix-valued Szegő polynomials; see, e.g., [4] for discussions on these polynomials, which are associated with matrix-valued measures on the unit circle. We are presently extending the approach of the present paper to investigate the conditioning of mappings analogous to (1.8) and (1.9) for matrix-valued Szegő polynomials in order to better understand the numerical properties of these polynomials.

We finally remark that the matrix (1.6) can be expressed as a product of  $n$  elementary unitary matrices,

$$H = G_1(\gamma_1)G_2(\gamma_2)\cdots G_{n-1}(\gamma_{n-1})\hat{G}_n(\tau),$$

where the  $G_j(\gamma_j)$ ,  $1 \leq j < n$ , are  $n \times n$  Givens matrices

$$G_j(\gamma_j) := \begin{bmatrix} I_{j-1} & & & \\ & -\gamma_j & \eta_j & \\ & \eta_j & \bar{\gamma}_j & \\ & & & I_{n-j-1} \end{bmatrix}, \quad \eta_j \in \mathbb{R}^+, \quad |\gamma_j|^2 + \eta_j^2 = 1,$$

and

$$\hat{G}_n(\tau) := \text{diag}[1, \dots, 1, -\tau].$$

Therefore small perturbations in  $H$  correspond to small perturbations in the  $\gamma_j$  and  $\tau$ . It follows that the mapping from  $[\omega, \lambda]$  to the coefficients  $[\gamma_1, \gamma_2, \dots, \gamma_{n-1}, \tau]$  is well conditioned when the mapping  $\mathcal{F}_n$  is well conditioned.

**2. The conditioning of  $\mathcal{F}_n$  and  $\tilde{\mathcal{F}}_n$ .** The aim of this section is to bound the condition number of the nonlinear mappings (1.8) and (1.9). Our investigation follows the approach used by Beckermann and Bourreau [3] in their study of the choice of modified moments with respect to an inner product on an interval on the real axis.

In order to define the condition number of a matrix-valued map  $\mathcal{A} : \mathbb{C}^m \ni x \mapsto A \in \mathbb{C}^{n \times n}$ , where  $x = [x_1, x_2, \dots, x_m]^t$ , we introduce

$$\frac{\partial A}{\partial x_j} := \lim_{\Delta x_j \rightarrow 0} \frac{\Delta A}{\Delta x_j}.$$

Here  $\Delta A$  is the perturbation of  $A$  caused by the error  $\Delta x_k$  in the data  $x_k$ . Then we could define the condition number of  $\mathcal{A}$  as

$$(2.1) \quad \kappa(\mathcal{A})(x) := \frac{\|x\|_2}{\|A\|_2} \left( \sum_{j=1}^m \left\| \frac{\partial A}{\partial x_j} \right\|_F^2 \right)^{1/2}.$$

Here and below  $\|\cdot\|_2$  denotes the Euclidean vector norm and  $\|\cdot\|_F$  the Frobenius norm. The condition number (2.1) has previously been used by Zhang [32, p. 573] in his analysis of the sensitivity of orthogonal polynomials of Sobolev-type to perturbations in the moments and modified moments. However, in order to obtain more meaningful results when the quantities involved are of different orders of magnitude, we allow nonnegative scaling factors  $\delta_j$ , similarly as Beckermann and Bourreau in their investigation [3]. Thus, introduce the scaled condition number

$$\kappa^{(\delta)}(\mathcal{A})(x) := \frac{1}{\|A\|_2} \left( \sum_{j=1}^m \left| \frac{x_j}{\delta_j} \right|^2 \right)^{1/2} \left( \sum_{j=1}^m \delta_j^2 \left\| \frac{\partial A}{\partial x_j} \right\|_F^2 \right)^{1/2},$$

which we minimize over all scaling factors  $\delta_j > 0$  to obtain

$$\hat{\kappa}(\mathcal{A})(x) := \inf_{\delta_j > 0} \kappa^{(\delta)}(\mathcal{A})(x) = \frac{1}{\|A\|_2} \sum_{j=1}^m |x_j| \left\| \frac{\partial A}{\partial x_j} \right\|_F.$$

It follows from the Cauchy inequality that  $\hat{\kappa}(\mathcal{A})(x) \leq \kappa^{(\delta)}(\mathcal{A})(x)$ . The corresponding condition number for the mapping  $\mathcal{F}_n$  is given by

$$(2.2) \quad \hat{\kappa}_{[\omega, \lambda]}(\mathcal{F}_n)([\omega, \lambda]) := \sum_{j=1}^n \left( \omega_j \left\| \frac{\partial H}{\partial \omega_j} \right\|_F + \left\| \frac{\partial H}{\partial \lambda_j} \right\|_F \right),$$

where we have used the fact that  $\|H\|_2 = 1$ . The subscript  $[\omega, \lambda]$  of  $\hat{\kappa}_{[\omega, \lambda]}$  indicates that we consider perturbations in all  $\omega_j$  and  $\lambda_j$ . We also are interested in the sensitivity of  $\mathcal{F}_n$  to perturbation in the  $\omega_j$ , only, and denote the the associated condition number by

$$(2.3) \quad \hat{\kappa}_{[\omega]}(\mathcal{F}_n)([\omega, \lambda]) := \sum_{j=1}^n \omega_j \left\| \frac{\partial H}{\partial \omega_j} \right\|_F.$$

Beckermann and Bourreau [3] studied the relation between Gauss quadrature data and the related symmetric tridiagonal Jacobi matrix. Using their techniques, we obtain upper bounds for the partial derivatives in (2.2). This yields the following bound for the condition number of  $\mathcal{F}_n$ .

**THEOREM 2.1.** *The condition number (2.2) satisfies the bound*

$$(2.4) \quad \hat{\kappa}_{[\omega, \lambda]}(\mathcal{F}_n)([\omega, \lambda]) \leq 7n + 6 \sum_{j=1}^n \omega_j \left( \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2},$$

where  $l'_k$  denotes the derivative of the Lagrange polynomial  $l_k$  defined by (1.12).

The proof of this bound requires a few auxiliary results, which we formulate as lemmas. The first of these recalls some decompositions of matrices related to the recursion coefficient matrix  $H$  derived by Gragg [12].

**LEMMA 2.2.** *Let*

$$V_n := \begin{bmatrix} 1 & 1 & \cdots & 1 \\ \lambda_1 & \lambda_2 & \cdots & \lambda_n \\ \vdots & \vdots & & \vdots \\ \lambda_1^{n-1} & \lambda_2^{n-1} & \cdots & \lambda_n^{n-1} \end{bmatrix},$$

$$\Lambda_n := \text{diag}[\lambda_1, \lambda_2, \dots, \lambda_n], \quad W_n := \text{diag}[\omega_1, \omega_2, \dots, \omega_n],$$

and  $S := V_n W_n$ . Then  $H^t$  has the spectral decomposition  $H^t = Q_n \Lambda_n Q_n^*$  with a unitary matrix  $Q_n$ . The columns are scaled to have positive first component. Thus, the  $j$ th column of  $Q_n$  has first component  $\omega_j > 0$ ; cf. (1.7). Moreover,  $S^*$  has the QR-factorization with unitary factor  $Q_n^*$  and an upper triangular factor with positive diagonal entries. Further,  $SS^* = \overline{M}_n$ , where  $M_n = [\mu_{j-k}]_{j,k=0}^{n-1}$  is the moment matrix.

A variant of Lemma 2.3 below is shown by Beckermann and Bourreau [3, Lemma 3] for the case when  $A \in \mathbb{R}^{n \times n}$  and  $x \in \mathbb{R}^n$ . The proof in [3] is based on results by Stewart [27, Theorem 3.1] on the perturbation of the QR-factorization of a real matrix. We modify [3, Lemma 3] in two ways: We allow  $A \in \mathbb{C}^{n \times n}$  and  $x \in \mathbb{C}^n$ , and only show a first order inequality in  $\|x\|_2$ . In our application of this result,  $\|x\|_2 = \mathcal{O}(\epsilon)$  with  $\epsilon$  “tiny.” In the following lemma and below, the symbols  $\doteq$  and  $\lesssim$  denote first order equalities and inequalities in  $\epsilon$ , respectively.

LEMMA 2.3. Consider  $A \in \mathbb{C}^{n \times n}$  and its perturbed counterpart  $\tilde{A} = (I_n + e_j x^t)A$ , with  $x \in \mathbb{C}^n$  and  $e_j$  the  $j$ th axis vector. Assume that  $\|x\|_2 = \mathcal{O}(\epsilon)$ . Then, for the QR-factorizations  $A = QR$ ,  $\tilde{A} = \tilde{Q}\tilde{R}$ , with  $R$  and  $\tilde{R}$  having real diagonal entries, we have

$$(2.5) \quad \|\tilde{Q} - Q\|_F \lesssim 3\|x\|_2.$$

*Proof.* The result can be shown by following the development of Stewart [27, pp. 516–517] for the special case when  $A = Q = R = I$ . This situation allows some simplifications, which yield (2.5). Stewart [27] shows his bounds for  $A \in \mathbb{R}^{n \times n}$  and  $x \in \mathbb{R}^n$ , but points out that his results can be extended to the complex case under the condition that the diagonal elements of the upper triangular factors in the QR-factorizations are real.  $\square$

We are in a position to show Theorem 2.1.

*Proof.* (i) For each fixed  $j \in \{1, 2, \dots, n\}$ , we consider the perturbed matrices  $\tilde{H}$ ,  $\tilde{Q}_n$ ,  $\tilde{\Lambda}_n$ , and  $\tilde{S}$  obtained by replacing  $\omega_j$  by  $\tilde{\omega}_j = \omega_j + \epsilon$  for some small  $|\epsilon|$ ,  $\tilde{\omega}_k = \omega_k$  for  $k \neq j$ , and  $\tilde{\lambda}_k = \lambda_k$  for  $k = 1, 2, \dots, n$ . Our aim is to estimate

$$\left\| \frac{\partial H}{\partial \omega_j} \right\|_F = \lim_{\epsilon \rightarrow 0} \frac{\|\tilde{H} - H\|_F}{|\epsilon|}.$$

Let  $A_j$  denote the  $j$ th row of the matrix  $A$ . Then

$$\tilde{S}^* = \tilde{W}_n V_n^* = \begin{bmatrix} \omega_1(V_n^*)_1 \\ \vdots \\ (1 + \frac{\epsilon}{\omega_j})\omega_j(V_n^*)_j \\ \vdots \\ \omega_n(V_n^*)_n \end{bmatrix} = (I_n + e_j x^t)S^*, \quad x := \left(\frac{\epsilon}{\omega_j}\right)e_j.$$

Application of Lemma 2.3 shows that

$$\|\tilde{Q}_n^* - Q_n^*\|_F \lesssim \frac{3|\epsilon|}{\omega_j}$$

for  $|\epsilon|$  sufficiently small. Thus, it follows from Lemma 2.2 that

$$\begin{aligned} \|\tilde{H} - H\|_F &= \|\tilde{H}^t - H^t\|_F = \|\tilde{Q}_n \Lambda_n \tilde{Q}_n^* - Q_n \Lambda_n Q_n^*\|_F \\ &= \|(\tilde{Q}_n - Q_n) \Lambda_n \tilde{Q}_n^* + Q_n \Lambda_n (\tilde{Q}_n^* - Q_n^*)\|_F \\ &\leq 2\|\tilde{Q}_n - Q_n\|_F \\ &\leq \frac{6|\epsilon|}{\omega_j}. \end{aligned}$$

Therefore,

$$\left\| \frac{\partial H}{\partial \omega_j} \right\|_F \leq \frac{6}{\omega_j}.$$

(ii) Similarly, for each fixed  $j \in \{1, 2, \dots, n\}$ , let  $\tilde{H}$ ,  $\tilde{Q}_n$ ,  $\tilde{\Lambda}_n$ , and  $\tilde{S}$  be the perturbed matrices obtained by replacing  $\lambda_j$  by  $\lambda_j + \epsilon$  for some small  $|\epsilon|$ ,  $\tilde{\lambda}_k = \lambda_k$  for  $k \neq j$ , and  $\tilde{\omega}_k = \omega_k$  for  $k = 1, 2, \dots, n$ . We would like to estimate

$$\left\| \frac{\partial H}{\partial \lambda_j} \right\|_F = \lim_{\epsilon \rightarrow 0} \frac{\|\tilde{H} - H\|_F}{|\epsilon|}.$$

It follows from

$$[0, 1, 2\lambda_j, \dots, (n-1)\lambda_j^{n-2}]^t = V_n^t [l'_1(\lambda_j), l'_2(\lambda_j), \dots, l'_n(\lambda_j)]^t$$

that

$$\begin{aligned} \tilde{S} &= \tilde{V}_n W_n \doteq V_n (I_n + \tilde{x} e_j^t) W_n, & \text{where } \tilde{x} &= \epsilon [l'_1(\lambda_j), \dots, l'_n(\lambda_j)]^t, \\ &= V_n W_n (I_n + x e_j^t), & \text{where } x &= \epsilon \left[ \frac{\omega_j}{\omega_1} l'_1(\lambda_j), \dots, \frac{\omega_j}{\omega_n} l'_n(\lambda_j) \right]^t. \end{aligned}$$

Application of Lemma 2.3 to  $\tilde{S}^* = (I_n + e_j x^*) W_n V_n^*$  with  $|\epsilon|$  sufficiently small yields

$$\|\tilde{Q}_n^* - Q_n^*\|_F \leq 3|\epsilon| \left( \omega_j^2 \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2}.$$

Therefore,

$$\begin{aligned} \|\tilde{H} - H\|_F &= \|\tilde{Q}_n \tilde{\Lambda}_n \tilde{Q}_n^* - Q_n \Lambda_n Q_n^*\|_F \\ &= \|(\tilde{Q}_n - Q_n) \tilde{\Lambda}_n \tilde{Q}_n^* + Q_n (\tilde{\Lambda}_n - \Lambda_n) \tilde{Q}_n^* + Q_n \Lambda_n (\tilde{Q}_n^* - Q_n^*)\|_F \\ &\leq 2\|\tilde{Q}_n - Q_n\|_F + \|\tilde{\Lambda}_n - \Lambda_n\|_F \\ &\leq 6|\epsilon| \left( \omega_j^2 \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2} + |\epsilon| \end{aligned}$$

and, consequently,

$$\left\| \frac{\partial H}{\partial \lambda_j} \right\|_F \leq 6 \left( \omega_j^2 \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2} + 1.$$

(iii) Substituting the bounds from (i) and (ii) into (2.2) completes the proof.  $\square$

The bound (2.4) indicates that the condition number (2.2) may be large if there are tiny  $\omega_j$  or if the nodes  $\lambda_j$  are clustered on a subarc of the unit circle. However, the above proof shows that  $\mathcal{F}_n$  is quite well-conditioned with respect to perturbations in the  $\omega_j$  only, even when some  $\omega_j > 0$  are tiny.

COROLLARY 2.4. *The condition number (2.3) satisfies the bound*

$$\hat{\kappa}_{[\omega]}(\mathcal{F}_n)([\omega, \lambda]) \leq 6n.$$

The following corollary simplifies the bound (2.4).

COROLLARY 2.5. *Let  $l_k$  be the Lagrange polynomials defined by (1.12) and introduce*

$$(2.6) \quad \alpha_k := \max_{|z|=1} |l_k(z)|, \quad 1 \leq k \leq n.$$

Then

$$(2.7) \quad \hat{\kappa}_{[\omega, \lambda]}(\mathcal{F}_n)([\omega, \lambda]) \leq 7n + 6n^2 \max_{1 \leq k \leq n} \frac{\alpha_k}{\omega_k}.$$

*Proof.* The Cauchy inequality yields

$$(2.8) \quad \sum_{j=1}^n \omega_j \left( \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2} \leq \left( \sum_{j,k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2}.$$

Since the Lagrange polynomials  $l_k$  are of degree  $n-1$ , Bernstein's inequality gives

$$\max_{|z|=1} |l'_k(z)| \leq (n-1) \max_{|z|=1} |l_k(z)|, \quad 1 \leq k \leq n.$$

Therefore,

$$\sum_{j,k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \leq n(n-1)^2 \sum_{k=1}^n \max_{|z|=1} \frac{|l_k(z)|^2}{\omega_k^2} \leq n^3 \sum_{k=1}^n \frac{\alpha_k^2}{\omega_k^2} \leq n^4 \max_{1 \leq k \leq n} \frac{\alpha_k^2}{\omega_k^2}.$$

Substitution into the right-hand side of (2.8) and using this bound in (2.4) shows (2.7).  $\square$

We turn to a bound for the condition number

$$(2.9) \quad \hat{\kappa}_{[\omega^2, \lambda]}(\tilde{\mathcal{F}}_n)([\omega^2, \lambda]) := \sum_{j=1}^n \left( \omega_j^2 \left\| \frac{\partial H}{\partial \omega_j^2} \right\|_F + \left\| \frac{\partial H}{\partial \lambda_j} \right\|_F \right)$$

for the mapping (1.9).

THEOREM 2.6. *The condition number (2.9) satisfies*

$$(2.10) \quad \hat{\kappa}_{[\omega^2, \lambda]}(\tilde{\mathcal{F}}_n)([\omega^2, \lambda]) \leq 4n + 6 \sum_{j=1}^n \omega_j \left( \sum_{k=1}^n \frac{|l'_k(\lambda_j)|^2}{\omega_k^2} \right)^{1/2},$$

where the  $l_k$  are given by (1.12).

*Proof.* The inequality is shown similarly as the bound (2.4). We use the same notation as in the proof of Theorem 2.1. Some details are omitted. Consider the

perturbed matrices  $\tilde{H}$ ,  $\tilde{Q}_n$ ,  $\tilde{\Lambda}_n$ , and  $\tilde{S}$  obtained by replacing  $\omega_j^2$  by  $\tilde{\omega}_j^2 = \omega_j^2 + \epsilon$  for some small  $|\epsilon|$ . Let  $\tilde{\omega}_k = \omega_k$  for  $k \neq j$ , and  $\tilde{\lambda}_k = \lambda_k$  for  $k = 1, 2, \dots, n$ . It follows from  $\tilde{\omega}_j \doteq \omega_j(1 + \frac{\epsilon}{2\omega_j^2})$  that

$$\tilde{S}^* \doteq (I_n + e_j x^*) S^*, \quad x := \left(\frac{\epsilon}{2\omega_j^2}\right) e_j.$$

Lemma 2.3 yields, for  $|\epsilon|$  sufficiently small,

$$\|\tilde{Q}_n^* - Q_n^*\|_F \leq \frac{3|\epsilon|}{2\omega_j^2}.$$

Therefore, by Lemma 2.2,

$$\|\tilde{H} - H\|_F \leq \frac{3|\epsilon|}{\omega_j^2}$$

and it follows that

$$\left\| \frac{\partial H}{\partial \omega_j} \right\|_F \leq \frac{3}{\omega_j^2}.$$

The bounds for the partial derivatives  $\left\| \frac{\partial H}{\partial \lambda_j} \right\|_F$  are the same as in the proof of Theorem 2.1. This shows the bound (2.10).  $\square$

**3. Numerical examples.** This section presents computed examples which illustrate the bounds of the previous section. The computations were carried in MATLAB with about 16 significant digits.

FIG. 3.1. *Example 3.1: The condition number of the map  $\mathcal{F}_n$  (dashed graph) and the bound of Theorem 2.1 (continuous graph).*

Example 3.1: Let the weights  $\omega_j$  and nodes  $\lambda_j$  be given by

$$\omega_1 := 1 - \frac{n-1}{n^3}, \quad \omega_j := \frac{1}{n^3}, \quad 2 \leq j \leq n,$$

$$\lambda_j := \exp(2\pi i(\frac{j-1}{n})), \quad 1 \leq j \leq n.$$

Figure 3.1 displays the condition number of the map  $\mathcal{F}_n$  and the upper bound of Theorem 2.1 for  $1 \leq n \leq 50$ . The quality of the bound can be seen to deteriorate as  $n$  increases. Nevertheless, the bound, and therefore also the condition number of  $\mathcal{F}_n$ , grow slower than exponentially with  $n$ .

Evaluation of the condition numbers of  $\mathcal{F}_n$  and  $\tilde{\mathcal{F}}_n$  shows the latter to be somewhat smaller than the former, in agreement with the fact that the bound of Theorem 2.6 is smaller than the bound of Theorem 2.1.  $\square$

FIG. 3.2. *Example 3.2: The condition number of the map  $\mathcal{F}_n$  (dashed graph) and the bound of Theorem 2.1 (continuous graph).*

TABLE 3.1  
*Example 3.2: Decrease of the smallest weights as a function of  $n$ .*

$n$	5	10	15	20	25	30
$\rho_n$	$3.8 \cdot 10^{-3}$	$6.9 \cdot 10^{-5}$	$3.4 \cdot 10^{-6}$	$2.7 \cdot 10^{-7}$	$2.9 \cdot 10^{-8}$	$4.0 \cdot 10^{-9}$

Example 3.2: Let the recursion coefficients be given by

$$(3.1) \quad \gamma_j := \frac{1}{\sqrt{j+1}}, \quad 1 \leq j < n.$$

Let  $\{\omega_j\}_{j=1}^n$  denote the weights of the  $n$ -point Szegő quadrature rule determined by the recursion coefficients (3.1) and  $\tau := -1$ . Define  $\rho_n := \min_{1 \leq j \leq n} \omega_j$ . Table 3.1 shows  $\rho_n$  to decrease rapidly as  $n$  increases. Therefore the condition number of the

map  $\mathcal{F}_n$  grows quite rapidly with  $n$ ; see Figure 3.2, which shows the condition number of  $\mathcal{F}_n$  as well as the upper bound of Theorem 2.1 for  $1 \leq n \leq 30$ . Note that the sum  $\sum_{j=1}^{n-1} \gamma_j^2$  grows unboundedly with  $n$ .  $\square$

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